

The spherical ensemble with external sources

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1 Outline

- ① Introduction
- ② Potential theory
- ③ Vector equilibrium problem
- ④ About the proof

1 Introduction: The spherical ensemble

n random points on the unit sphere $\mathbb{S}^2 \subset \mathbb{R}^3$ with joint probability density

$$\frac{1}{Z_n} \prod_{i=1}^{n-1} \prod_{j=i+1}^n \|x_i - x_j\|^2$$

- ▶ Determinantal point process
- ▶ Eigenvalues of a random matrix Krishnapur (2009)

Take independent Ginibre matrices A and B .

Compute eigenvalues of AB^{-1} .

Put points on \mathbb{S}^2 with inverse stereographic projection.

1 The spherical ensemble with external charges

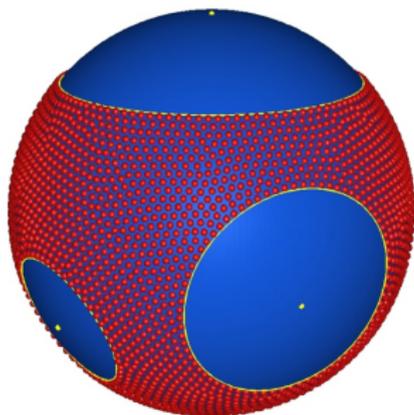
Fix $p_0, \dots, p_r \in \mathbb{S}^2$ and $a_0, \dots, a_r > 0$. It creates probability density for points x_1, \dots, x_n on the sphere

$$\frac{1}{Z_n} \prod_{i=1}^{n-1} \prod_{j=i+1}^n \|x_i - x_j\|^2 \cdot \prod_{j=1}^n \prod_{k=0}^r \|x_j - p_k\|^{2na_k}$$

Each p_k creates a region around itself that is free from x_j 's in large n limit.

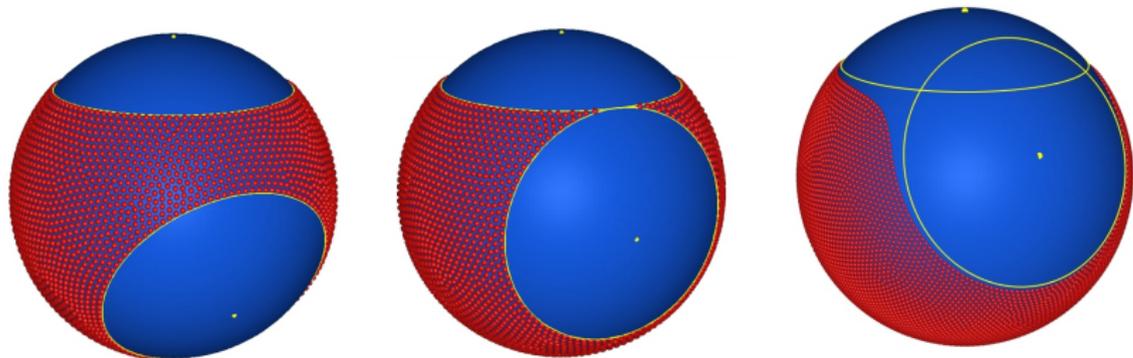
When a_k 's are small these regions are spherical caps

Brauchart, Dragnev, Saff,
Womersley



1 The spherical ensemble with external charges

When charges a_k are larger the regions start to overlap.



Figures from BDSW

- ▶ What region D is occupied by the free charges?
- ▶ D is known as the **droplet**.

1 Normal matrix model

Eigenvalues of **complex Ginibre matrix** have joint density

$$\frac{1}{Z_n} \prod_{i=1}^n \prod_{j=i+1}^n |z_i - z_j|^2 \prod_{j=1}^n e^{-n|z_j|^2}$$

- ▶ Eigenvalues fill out a disk as $n \rightarrow \infty$.

Ginibre model with **external sources** (also arises as eigenvalues in a normal matrix model)

$$\frac{1}{Z_n} \prod_{i=1}^n \prod_{j=i+1}^n |z_i - z_j|^2 \prod_{j=1}^n e^{-n|z_j|^2} \prod_{j=1}^n \prod_{k=0}^r |z_j - q_k|^{2na_k}$$

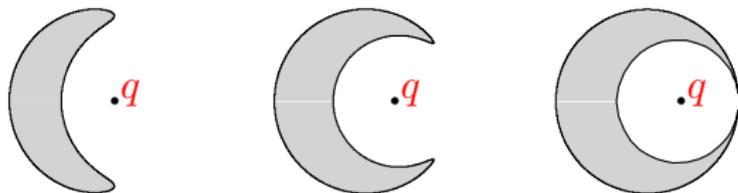
- ▶ Studied in detail for case $r = 0$ by

Balogh, Bertola, Lee, McLaughlin 2015

1 Droplet

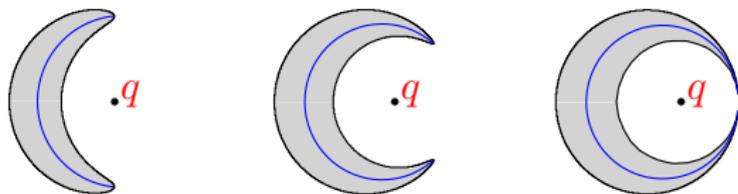
$$\frac{1}{Z_n} \prod_{i=1}^n \prod_{j=i+1}^n |z_i - z_j|^2 \prod_{j=1}^n e^{-n|z_j|^2} \prod_{j=1}^n |z_j - q|^{2na}$$

Change in **droplet** as a increases.



1 Motherbody

Average characteristic polynomial $P_n(z) = \mathbb{E} \left[\prod_{j=1}^n (z - z_j) \right]$ is **planar orthogonal polynomial** with zeros that cluster around a contour as $n \rightarrow \infty$, known as the **motherbody**



- ▶ Planar orthogonality can be rewritten as orthogonality along a closed contour around q .

Balogh, Bertola, Lee, McLaughlin 2015, Lee, Yang 2019

- ▶ This allows a Riemann-Hilbert asymptotic analysis

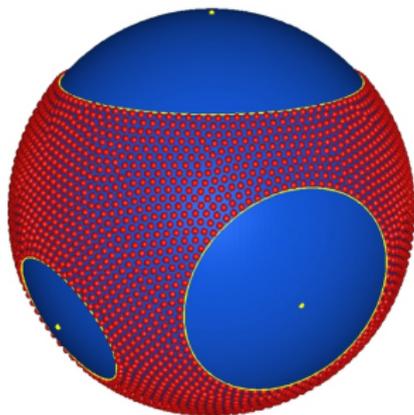
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2 Potential theory

Back to spherical model

$$\frac{1}{Z_n} \prod_{i=1}^{n-1} \prod_{j=i+1}^n \|x_i - x_j\|^2 \cdot \prod_{j=1}^n \prod_{k=0}^r \|x_j - p_k\|^{2i}$$



Fixed charge distribution $\sigma = \sum_{k=0}^r a_k \delta_{p_k}$

Limiting distribution μ_σ of free charges is unique probability measure on \mathbb{S}^2 that satisfies

$$U^{\mu_\sigma} + U^\sigma = \ell \quad \text{on } D = \text{supp}(\mu_\sigma),$$

$$U^{\mu_\sigma} + U^\sigma \geq \ell \quad \text{on } \mathbb{S}^2.$$

Notation: $U^\mu(x) = \int \log \frac{1}{\|x - y\|} d\mu(y)$

2 Potential theory

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Limiting distribution μ_σ of free charges is unique probability measure on \mathbb{S}^2 that satisfies

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Lemma

$$\mu_\sigma = (\lambda(D))^{-1} \lambda_D$$

where λ_D is the restriction to D of the normalized Lebesgue measure λ on \mathbb{S}^2 .

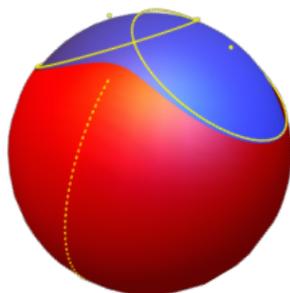
2 The motherbody

A **motherbody** for D is a probability measure σ^* on a one-dimensional subset of D such that

$$\begin{aligned}U\sigma^* &= (\lambda(D))^{-1}U\lambda_D + \ell^* && \text{on } \mathbb{S}^2 \setminus D, \\U\sigma^* &\geq (\lambda(D))^{-1}U\lambda_D + \ell^* && \text{on } D.\end{aligned}$$

In case of two points p_0, p_1 with sufficiently large equal charges $a_0 = a_1$, $\mathbb{S}^2 \setminus D$ is an **ellipse** after stereographic projection onto the complex plane.

The motherbody is minimizer of **equilibrium problem with external field**
Criado del Rey - K



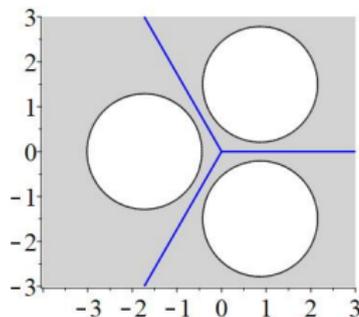
2 More external charges

We generalize this to $r + 1$ points p_0, \dots, p_r **with symmetry**

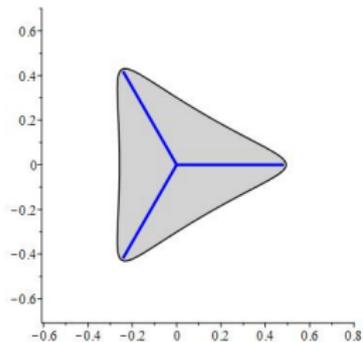
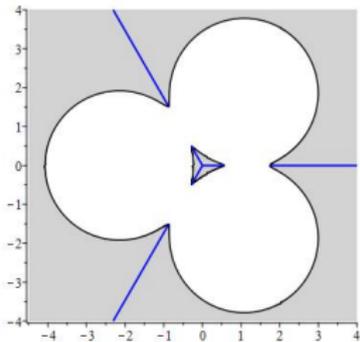
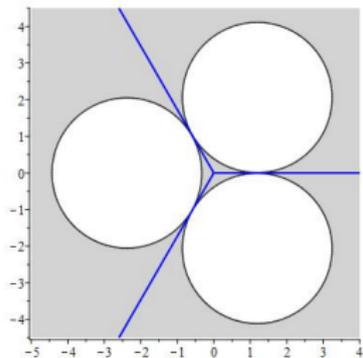
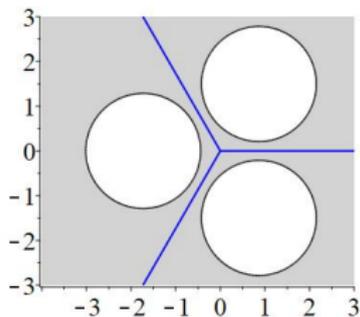
- ▶ p_j 's are located symmetrically around the north pole
- ▶ Equal charges $a_0 = a_1 = \dots = a_r = a > 0$

Motherbody will be supported on $r + 1$ **meridians** on the sphere, connecting the north and south poles.

After stereographic projection these are $r + 1$ **half-rays** in the complex plane.



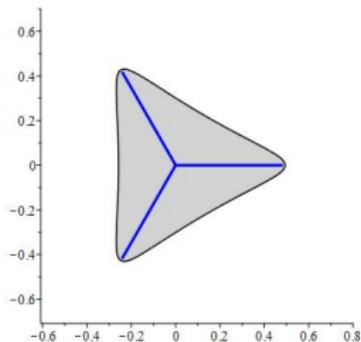
2 More pictures for $r = 2$ (three charges)



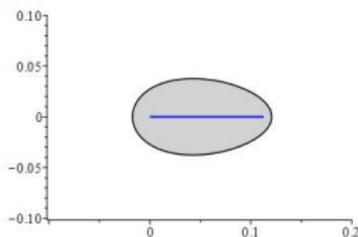
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3 Remove the symmetry



$$z \mapsto z^{r+1}$$



We expect **three cases** for the support Σ of the motherbody (after removing the symmetry)

- ▶ **Bounded interval support (BIS):** $\Sigma = [0, x_1]$
- ▶ **Two interval support (TIS):** $\Sigma = [0, x_1] \cup [x_2, \infty)$
- ▶ **Full interval support (FIS):** $\Sigma = [0, \infty)$

3 Vector equilibrium problem with r measures

Energy functional with input $q > 0$, $0 < t < 1$ and integer $r \geq 1$

$$\begin{aligned} \mathcal{E}(\mu_1, \dots, \mu_r) &= \sum_{j=1}^r I(\mu_j) - \sum_{j=1}^{r-1} I(\mu_j, \mu_{j+1}) \\ &\quad - \frac{1-t}{t} \int \log |x + q^{-1}| d\mu_1(x) + \frac{r+t}{t} \int \log |x - (-1)^r q| d\mu_r(x) \end{aligned}$$

Notation:

$$\begin{aligned} I(\mu, \nu) &= \iint \log \frac{1}{|x-y|} d\mu(x) d\nu(y) \\ I(\mu) &= I(\mu, \mu) \end{aligned}$$

3 Vector equilibrium problem with r measures

Energy functional with input $q > 0$, $0 < t < 1$ and integer $r \geq 1$

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Minimize $\mathcal{E}(\mu_1, \dots, \mu_r)$ under conditions

- ▶ μ_j is supported on $[0, \infty)$ if j is odd, and on $(-\infty, 0]$ if j is even,

- ▶ $\int d\mu_j = 1 + \frac{j-1}{t}$ for $j = 1, \dots, r$.

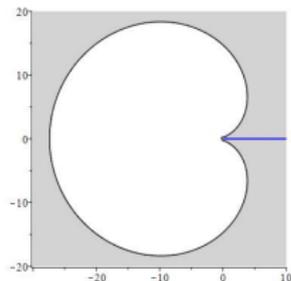
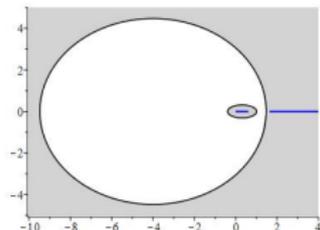
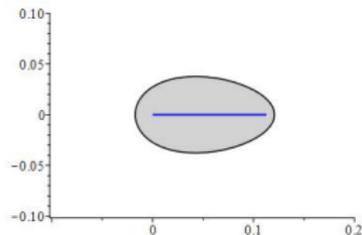
3 Main result

Theorem

- (a) *The vector equilibrium problem has a **unique minimizer**.*
- (b) *The first component μ_1 has **support** Σ_1 that is either $[0, x_1]$, or $[0, x_1] \cup [x_2, \infty)$ or $[0, \infty)$.*
- (c) *$\Phi(z) = t \int \frac{d\mu_1(s)}{z-s} + \frac{1-t}{z+q^{-1}}$ is an **algebraic function with meromorphic continuation to an $r+1$ sheeted Riemann surface***
- (d) *There is a **domain** U containing Σ_1 with boundary*

$$\partial U : \quad z\Phi(z) = \frac{|z|^{\frac{2}{r+1}}}{1 + |z|^{\frac{2}{r+1}}}$$

3 Pictures (on different scales)



BIS: closed contour **TIS:** two contours **FIS:** closed contour

∂U is boundary of **shaded region** U

- ▶ **BIS:** $\Sigma_1 = [0, x_1]$ and U is bounded
- ▶ **TIS:** $\Sigma_1 = [0, x_1] \cup [x_2, \infty)$ and U has one bounded and one unbounded component
- ▶ **FIS:** $\Sigma_1 = [0, \infty)$ and U is unbounded

3 Main result (cont.)

Theorem

(e) $\Omega = \{z \in \mathbb{C} \mid z^{r+1} \in U\}$ is mapped by inverse stereographic projection to the **droplet** D

The points p_0, \dots, p_r correspond to **equidistant points** in the plane on a circle with **radius** $q^{-\frac{1}{r+1}} > 1$

The **charges** are

$$a = \frac{1-t}{t(r+1)} > 0$$

(f) The measure μ_1 gives the **motherbody** of D (after re-introducing the $r+1$ -fold symmetry and mapping back to the unit sphere)

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4 Vector equilibrium problem

- ▶ Vector equilibrium problem is **weakly admissible**
- ▶ The measures μ_2, \dots, μ_r have full supports $(-\infty, 0]$ or $[0, \infty)$ as they are **balayage measures**

$$2\mu_k = \text{Bal}(\mu_{k-1} + \mu_{k+1}; (-1)^{k+1}[0, \infty)), \quad \text{for } k = 2, 3, \dots, r$$

4 Vector equilibrium problem

- ▶ The measures μ_2, \dots, μ_r have full supports $(-\infty, 0]$ or $[0, \infty)$ as they are **balayage measures**
- ▶ μ_1 also minimizes

$$\begin{aligned} & \frac{1}{2} \iint \log \frac{1}{|x-y|} d\mu(x) d\mu(y) \\ & \quad + \frac{1}{2} \iint \log \frac{1}{|x^{1/r} - y^{1/r}|} d\mu(x) d\mu(y) \\ & + \int \left(\left(1 - \frac{1}{t}\right) \log \left(x + q^{-1}\right) + \left(1 + \frac{r}{t}\right) \log \left(x^{1/r} + q^{1/r}\right) \right) d\mu(x) \end{aligned}$$

among probability measures μ on $[0, \infty)$

Equilibrium problem for **Muttalib-Borodin** ensemble

Claeys Romano 2014, K 2016, Molag 2020

4 The support of μ_1

μ_1 is the balayage measure $\text{Bal}\left(-\left(\frac{1}{t}-1\right)\delta_{-q} + \mu_2; [0, \infty)\right)$
provided that this is **positive**

- ▶ This happens in **FIS** (full interval support) case.

There are four possibilities for the **support** Σ_1 of μ_1

- ▶ $\Sigma_1 = [0, x_1]$ with $0 < x_1 < \infty$ **BIS**
- ▶ $\Sigma_1 = [0, x_1] \cup [x_2, \infty)$ with $0 < x_1 < x_2 < \infty$ **TIS**
- ▶ $\Sigma_1 = [0, \infty)$ **FIS**
- ▶ $\Sigma_1 = [x_2, \infty)$ with $0 < x_2 < \infty$ **UIS**

UIS (unbounded interval support) does not happen for $0 < q \leq 1$, **BIS** does not happen for $q \geq 1$.

This is proved with **iterated balayage** algorithm

Dragnev Kuijlaars 1999

4 Monotonicity properties

$\mu_1 = \mu_{1,t}$ and $\Sigma_1 = \Sigma_{1,t}$ vary with $t \in (0, 1)$

Fix $0 < q < 1$

Theorem

- (a) $\Sigma_{1,t}$ *increases with t .*
- (b) $t\mu_{1,t}$ *increases with t .*
- (c) **There are two critical values** $0 < t_{1,cr} < t_{2,cr} < 1$ **such that**
 - BIS** in case $0 < t \leq t_{1,cr}$
 - TIS** in case $t_{1,cr} < t < t_{2,cr}$
 - FIS** in case $t_{2,cr} \leq t < 1$

(a) is implied by (b) and (b) follows from further analysis of iterated balayage algorithm (+some other tricks)

(c) follows from explicit calculations in BIS and FIS cases.

4 Riemann surface

Riemann surface \mathcal{R} with $r + 1$ sheets

$$(\Sigma_j = \text{supp}(\mu_j))$$

$$\mathcal{R}^{(1)} = \mathbb{C} \setminus \Sigma_1,$$

$$\mathcal{R}^{(j)} = \mathbb{C} \setminus (\Sigma_{j-1} \cup \Sigma_j) \quad \text{for } j = 2, \dots, r,$$

$$\mathcal{R}^{(r+1)} = \mathbb{C} \setminus \Sigma_r.$$

► **Meromorphic function** Φ is given on j th sheet by

$$\Phi^{(j)}(z) = t \int \frac{d\mu_j(s)}{z-s} - t \int \frac{d\mu_{j-1}(s)}{z-s}, \quad z \in \mathcal{R}^{(j)}$$

with $\mu_0 = \left(1 - \frac{1}{t}\right) \delta_{-q^{-1}}$ and $\mu_{r+1} = \left(1 + \frac{r}{t}\right) \delta_{(-1)^r q}$

► $z\Phi(z)$ has degree 2.

4 Droplet and spherical Schwarz function

▶ $\partial U : z\Phi(z) = \frac{|z|^{\frac{2}{r+1}}}{1 + |z|^{\frac{2}{r+1}}}$ is boundary of domain U .

▶ If $\Omega = \{z \mid z^{r+1} \in U\}$ and $S(z) = z^r \Phi(z^{r+1})$ then

$$\partial\Omega : S(z) = \frac{\bar{z}}{1 + |z|^2}$$

$S(z)$ is the **spherical Schwarz function** of $\partial\Omega$

▶ **Monotonicity** of Ω in parameter t is crucial to prove that Ω is the stereographic projection of the **droplet**.

4 Final remark

Thank you for your attention !